

# Kinetic aspects of the vortex-induced reconnection in collisionless plasmas: 2D & 3D full PIC simulations

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Simulation codes:

Electromagnetic relativistic particle-in-cell code developed in Japan

Electromagnetic relativistic particle-in-cell code developed in LANL (VPIC)

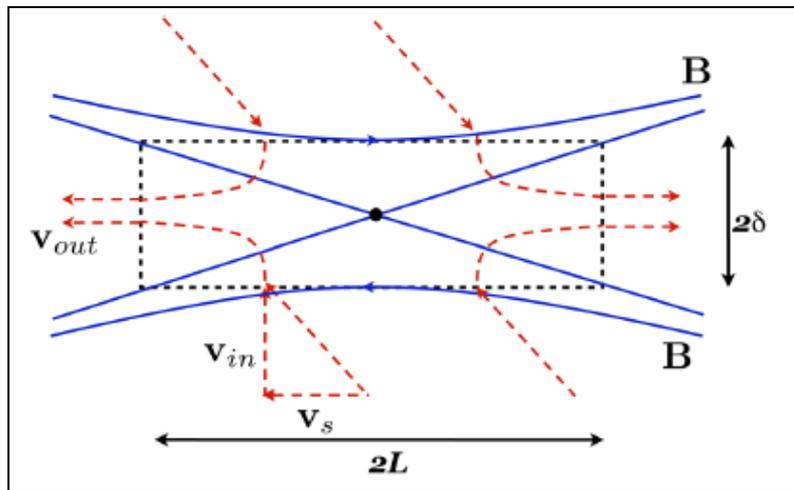
Computers:

LANL's and JAXA's supercomputers for 2D

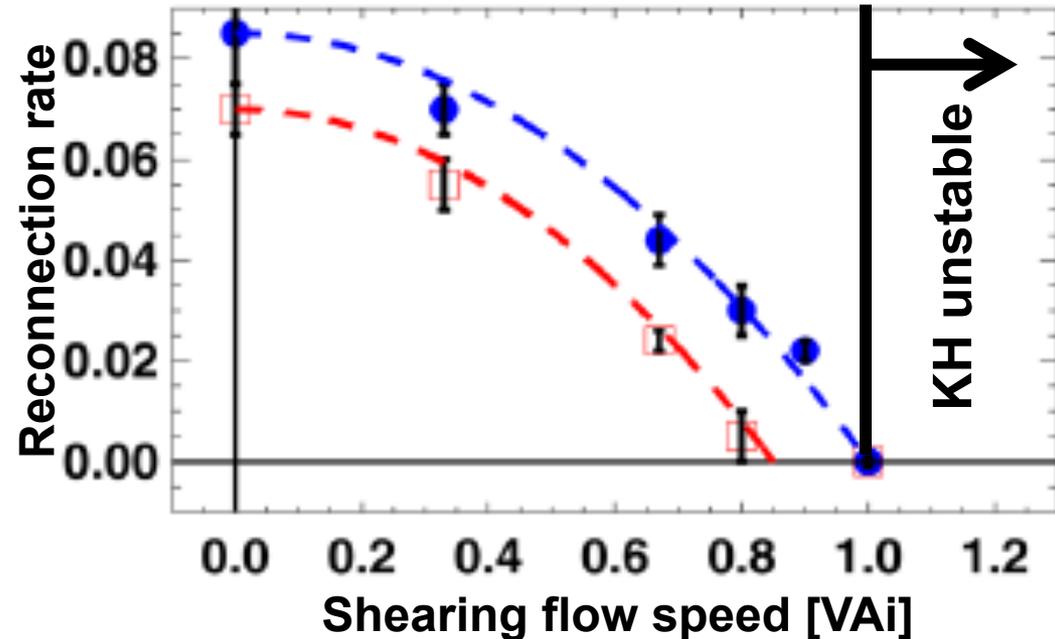
Jaguar (Oak Ridge Leadership Computing Facility) for 3D



# Reconnection + velocity shear



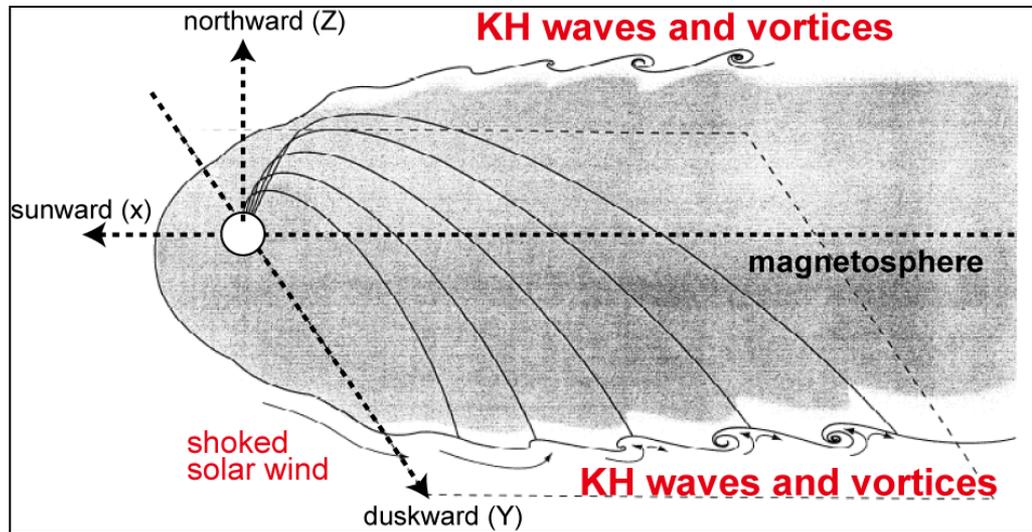
Reconnection region with velocity shear



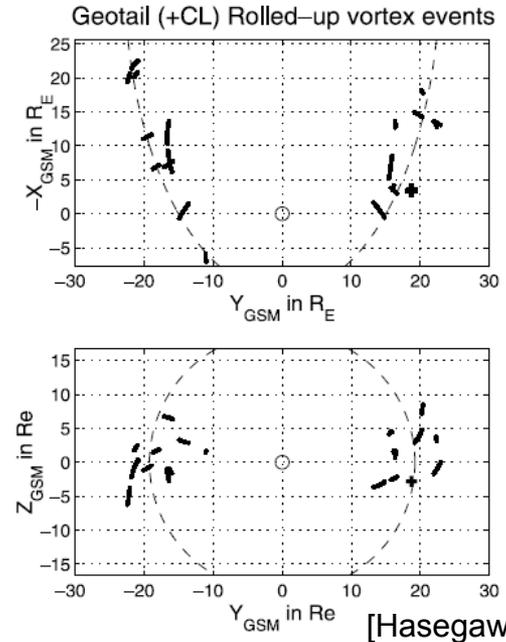
[Cassak & Otto, 2011]

- ◆ Basically, the velocity shear reduces the reconnection rate [e.g., Cassak & Otto, 2011].
- ◆ BUT, super Alfvénic shearing flow produces the Kelvin-Helmholtz instability.  
→ The resulting KH vortex can induce reconnection.

# KH vortex at the Earth's magnetopause



[Fairfield et al., 2000]



[Hasegawa et al., 2006]

◆ At the planetary magnetopause, the velocity shear between solar wind and magnetospheric plasmas increases with down tail distance.

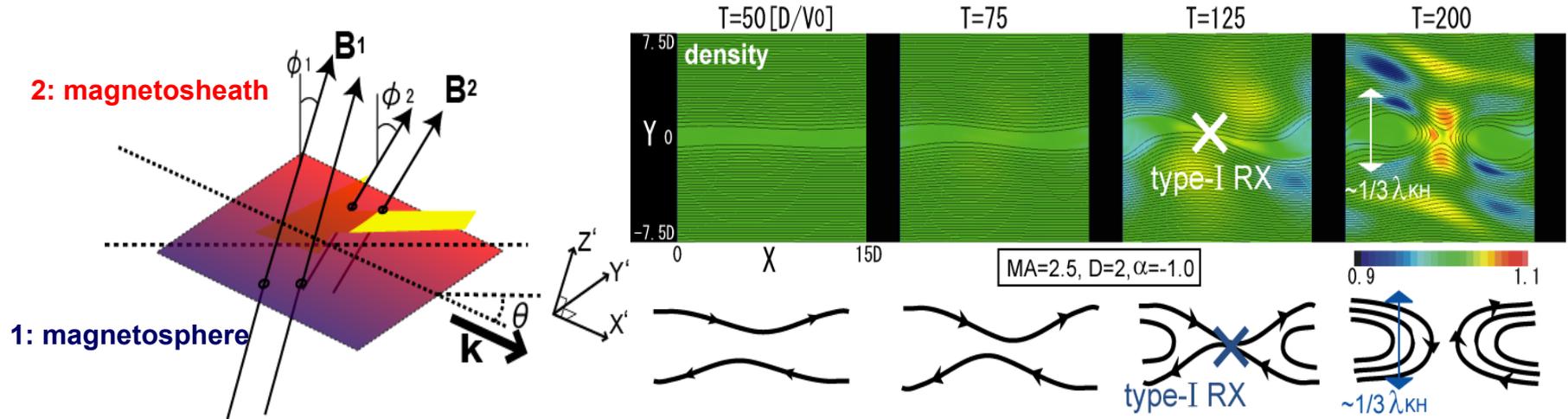
→ KH vortices would grow along the tail-flank magnetopause.

→ Indeed, in-situ observations have confirmed the rolled-up vortices there.

[e.g., Fairfield et al., 2000; Hasegawa et al., 2004, 2006, using Geotail, Cluster, THEMIS].

**Q. Can the vortex-induced reconnection also occur there?**

# VIR at the Earth's magnetopause



The type-I VIR by two-fluid simulation [Nakamura et al., 2008]

## A. Yes.

→ Linear analyses suggested that **the vortex-induced reconnection (VIR) commonly appear** at the Earth's magnetopause [Nakamura et al., 2006, 2008].

← This is because the magnetic shear (the current sheet) always exists at the magnetopause.

◆ Note that **the VIR can lead to plasma mixing** along reconnected field lines, and would form the part of the tail-flank low-latitude boundary layer (LLBL).

# **1. 2D symmetric case** [Nakamura et al., JGR, 2011]

(Plasmas and fields are symmetric across the boundary)

# **2. 2D asymmetric case** [Nakamura et al., in prep.]

(considering **density jump** across the boundary)

# **3. 3D symmetric case** [On-going]

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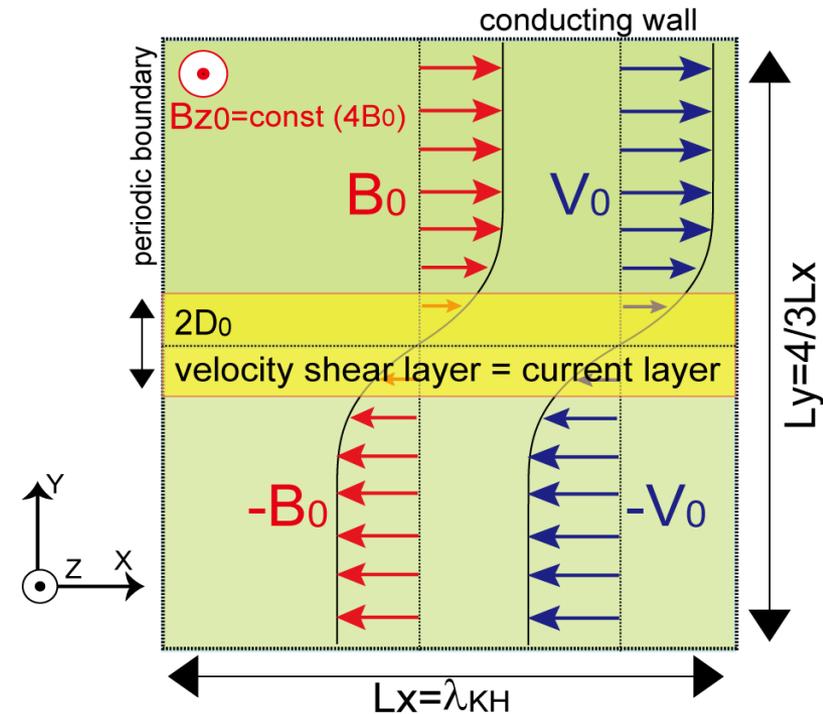
**3. 3D symmetric case** [On-going]

# Symmetric case [Nakamura et al., JGR, 2011]

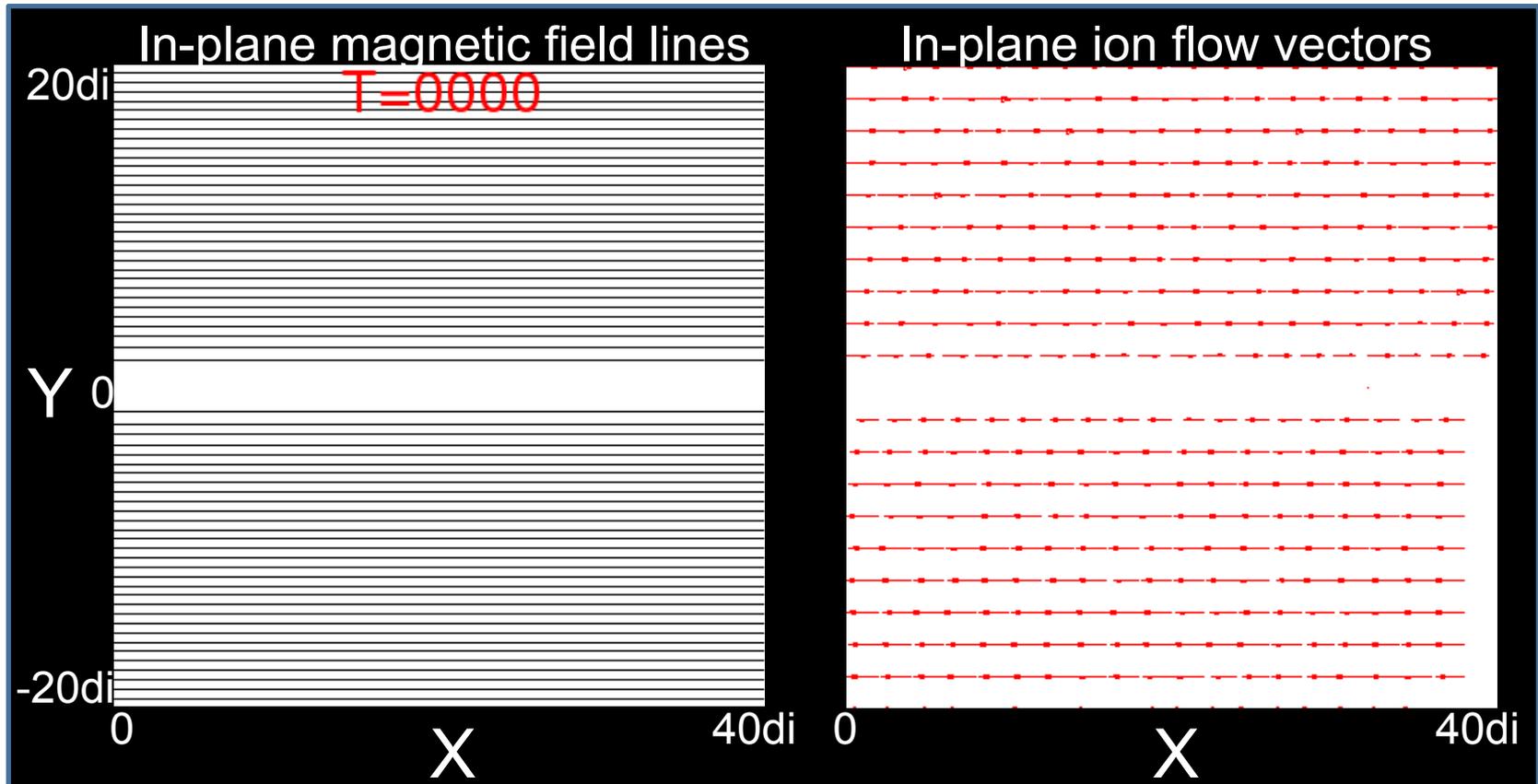
Method : 2.5-dimensional full kinetic PIC simulation developed in Japan [Hoshino, 1987]

## -Initial parameters-

- $N_0 = N_{i0}$  (uniform)
- $B_{z0} = 4 \cdot B_0$  (uniform)
- $B_{x0} = B_0 \cdot \tanh(Y/D_0)$  (current sheet)
- $V_{x0} = V_0 \cdot \tanh(Y/D_0)$  (velocity shear layer)
- $D_0 = 2.0$  [di] (MHD-scale boundary layer)
- $MA = V_0/V_{A0} = 4.375$  (strong KHI)  
(KHI can grow when  $MA > 2$  [Miura & Pritchett, 1982] )
- $L_x = \lambda_{KH} = 20D_0$   
(~fastest growing KH mode [Miura & Pritchett, 1982] )
- Other parameters:
  - $T_{i\_cs}/T_{e\_cs} = 1/8$ ,  $T_{i\_bg}/T_{e\_bg} = 1$
  - $M_i/M_e = 25, 100$
  - $\omega_{pe}/\Omega_e = 2.0$
  - 100 particles/grid

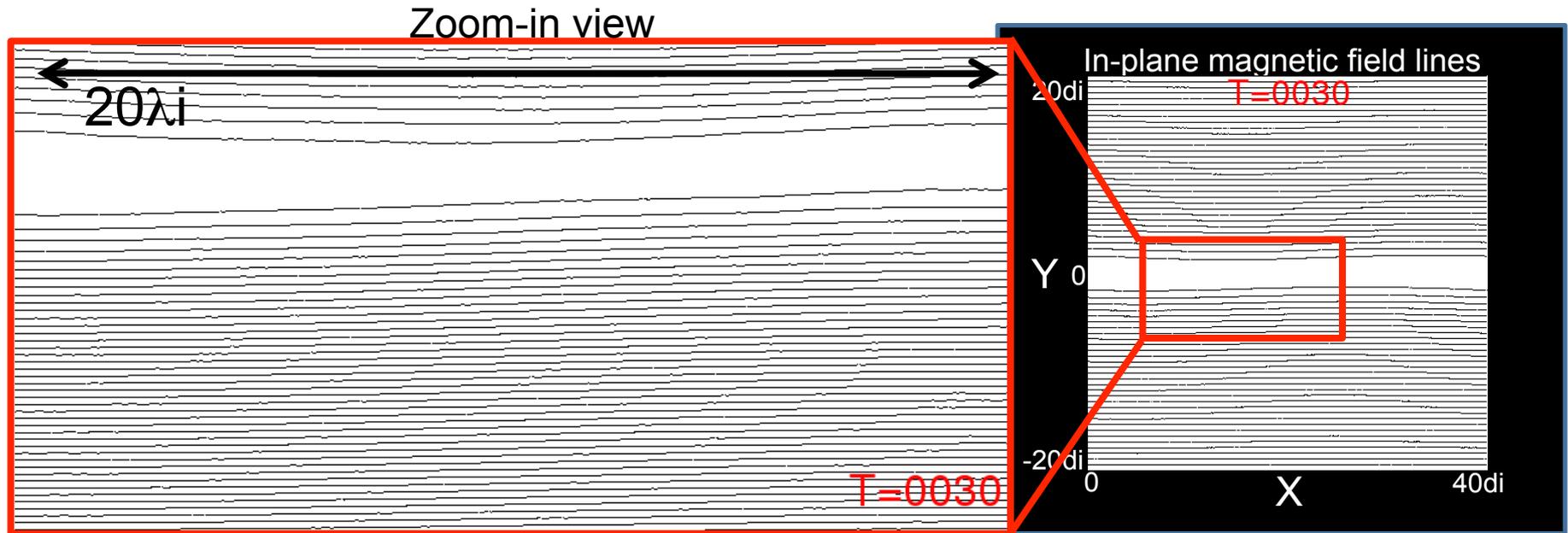


# Overview of symmetric case [Nakamura et al., 2011]



- ◆ (T<50) The KHI grows and locally compresses the current sheet.
- ◆ (T~50) Multiple reconnection occurs at the compressed thin current sheet.
- ◆ (T~80) Finally, the KH vortex is highly rolled-up as a large magnetic island.

# Multiple island formation [Nakamura et al., 2011]



(T=30-50)

- ◆ The strong vortex flow produces a thin and long current sheet.

(T=60-80)

- ◆ Multiple magnetic islands are formed and move toward the main body of the vortex.

(T~80)

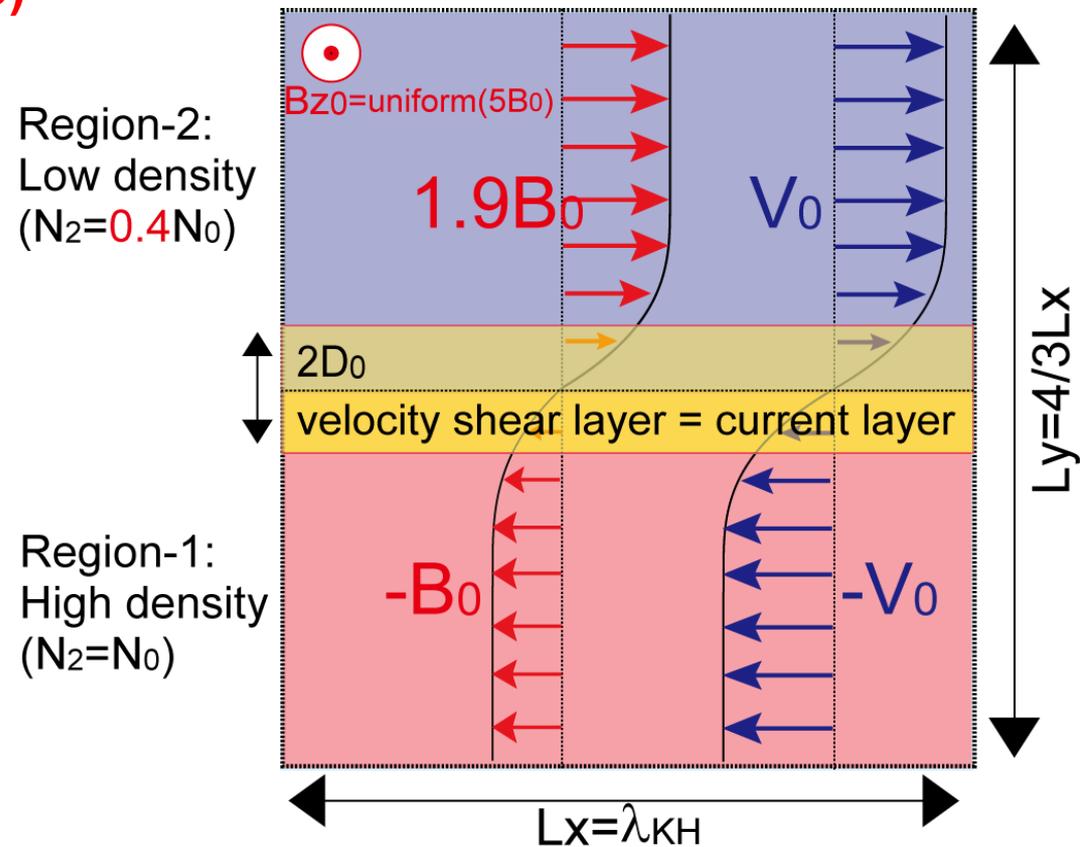
- ◆ Multiple magnetic islands are incorporated in turn into the vortex body via the re-connection.

# Asymmetric case [Nakamura et al., in prep.]

Method: 2.5-dimensional full kinetic PIC simulation developed in LANL (VPIC) [Bowers., 2008,2009]

## -Initial parameters-

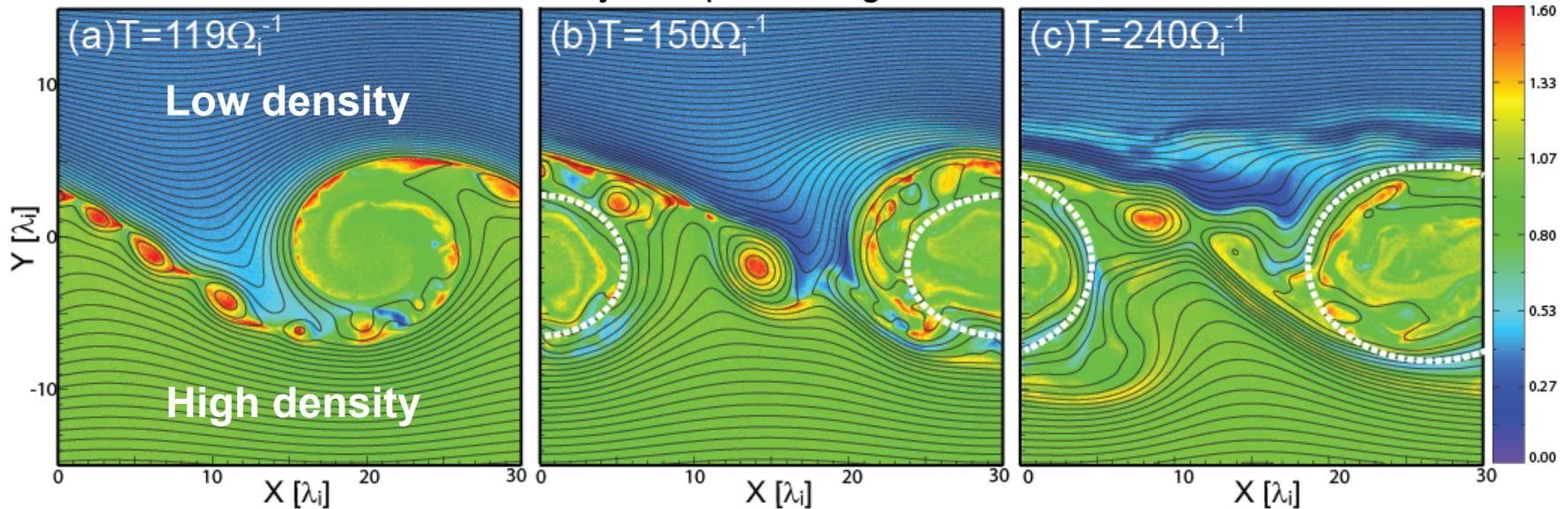
- $N_1 : N_2 = 1 : 0.4$  (density jump)
- $B_{x1} : B_{x2} = 1 : 1.875$
- $B_{z0} = 5 \cdot B_0$  (uniform)
- $V_{x0} = V_0 \cdot \tanh(Y/D_0)$
- $D_0 = 2.0$  [di] (MHD-scale)
- $MA = V_0 / V_{Ax1} = 7.0$  (strong KHI)
- $L_x = \lambda_{KH} = 15D_0$
- Other parameters:
  - $T_i / T_e = 2.5$
  - $M_i / M_e = 25, 100$
  - $\omega_{pe} / \Omega_e = 1.25$
  - 100 particles/grid



# Overview of asymmetric case

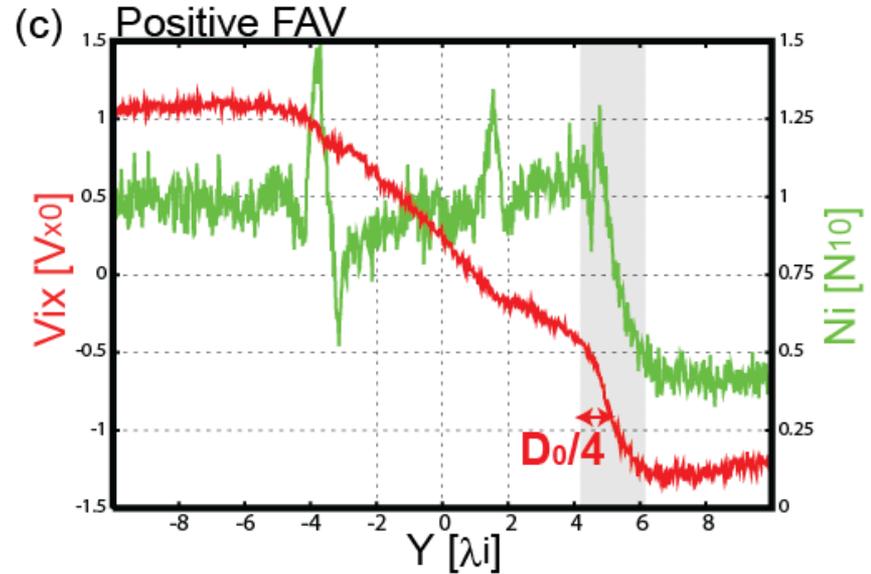
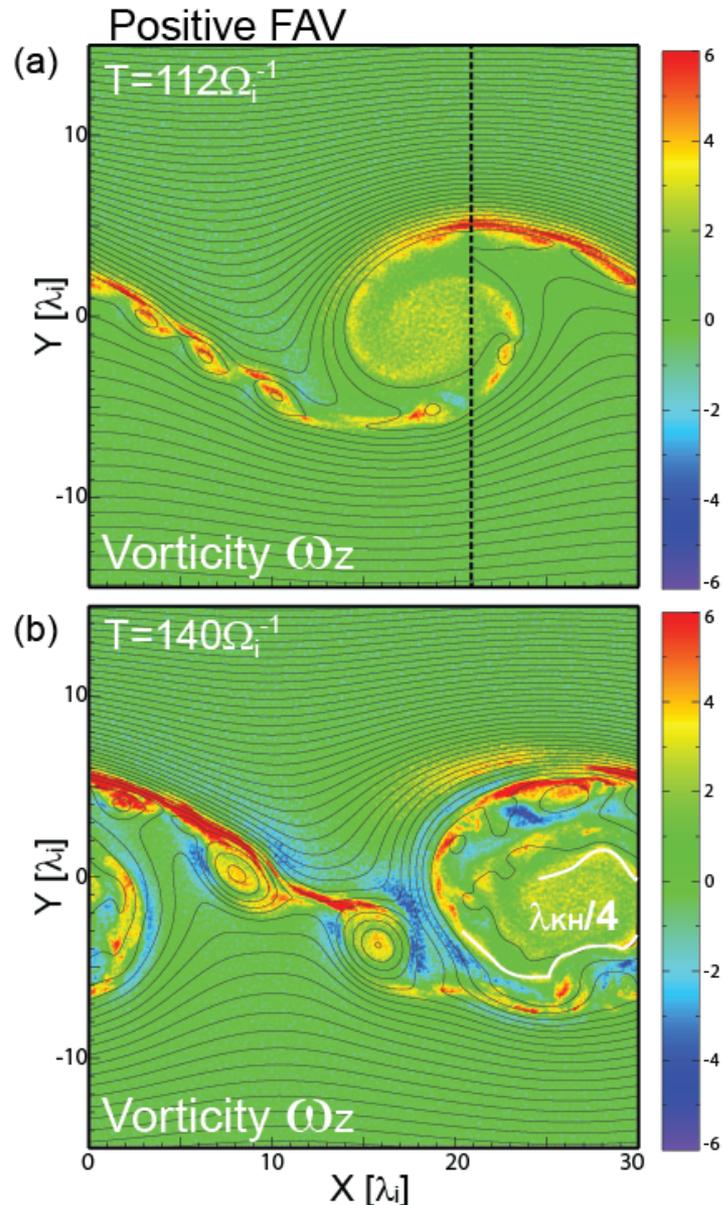
$$V_0=7.0V_{iA1}, D_0=2.0\lambda_i, N_1 : N_2=1 : 0.4$$

Ion density & in-plane magnetic field lines



- ◆ ( $T \sim 119$ ) Multiple islands are formed in the compressed current sheet.
- ◆ ( $T \sim 150$ ) **Secondary KH waves** grow within the parent vortex.
- ◆ ( $T \sim 240$ ) **Turbulence** is produced within the vortex.

# Unstable secondary shear layer

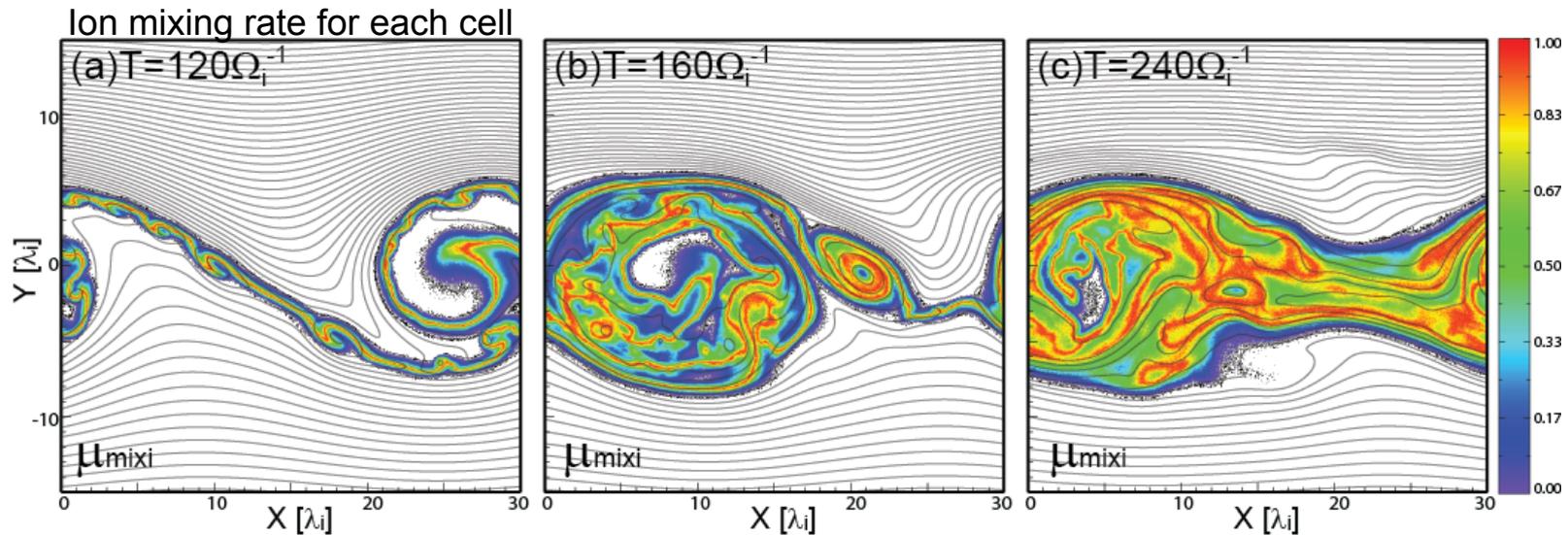


◆ The density jump leads to the formation of the **secondary velocity shear layer**.

◆ The VIR releases in-plane magnetic field within the vortex.

◆ The **secondary shear layer is always unstable for the secondary KHI within the vortex**.

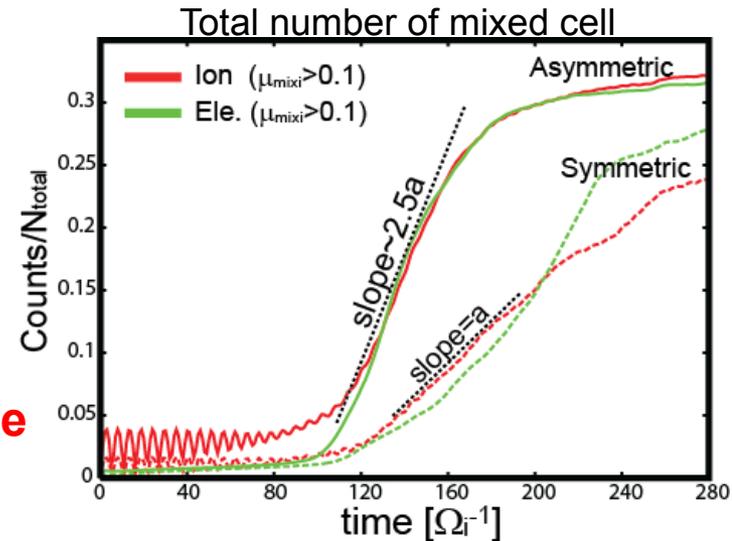
# Plasma mixing in asymmetric case



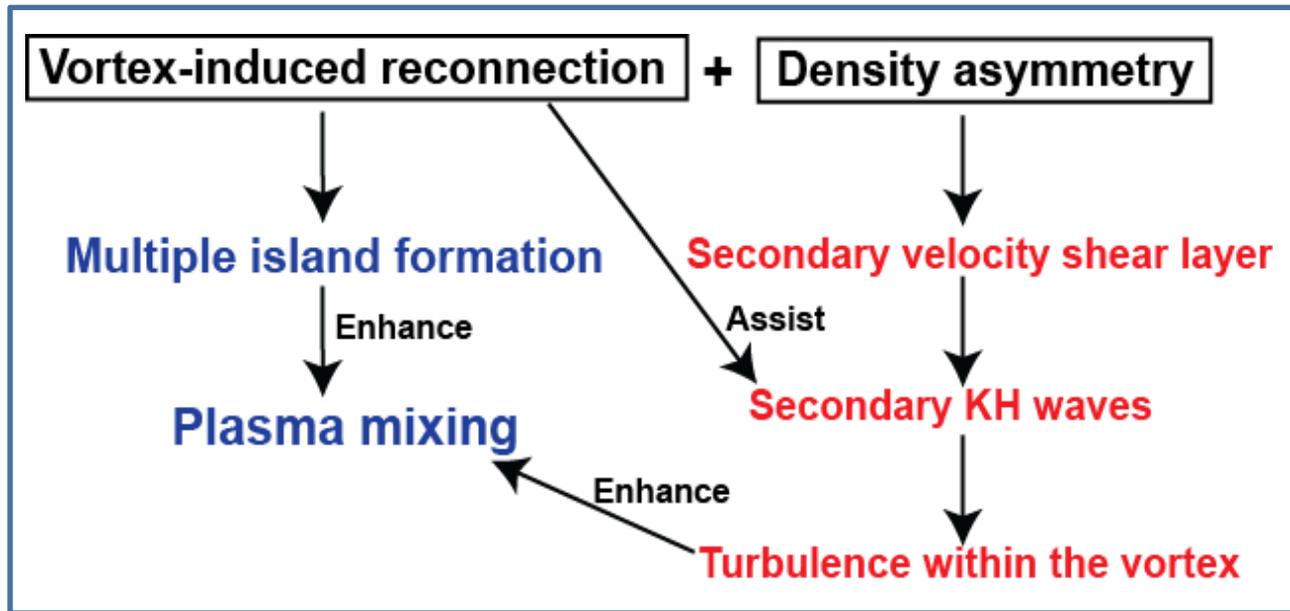
◆ ( $T < 120$ ) The vortex-induced reconnection (VIR) starts causing plasma mixing as well as the symmetric case.

◆ ( $T > 120$ ) Plasma mixing is further enhanced by **2ndary KH waves** and resulting **turbulence** unlike the symmetric case.

◆ **Consequently, plasma mixing in asymmetric case progresses more quickly than symmetric case.**



# Summary (2D)



- The **VIR** process accompanies the **multiple island formation** in the compressed current sheet.
  - The multiple island formation leads to the **efficient plasma mixing** within the vortex.
- The **density jump** across the shear layer leads to the formation of the **secondary shear layer**.
- The VIR can release the in-plane magnetic field within the vortex, and hence **the secondary KH instability** is always unstable within the vortex.
- Resulting secondary KH waves eventually cause **turbulence within the vortex**.
  - The turbulence formation **enhances the plasma mixing** within the vortex.

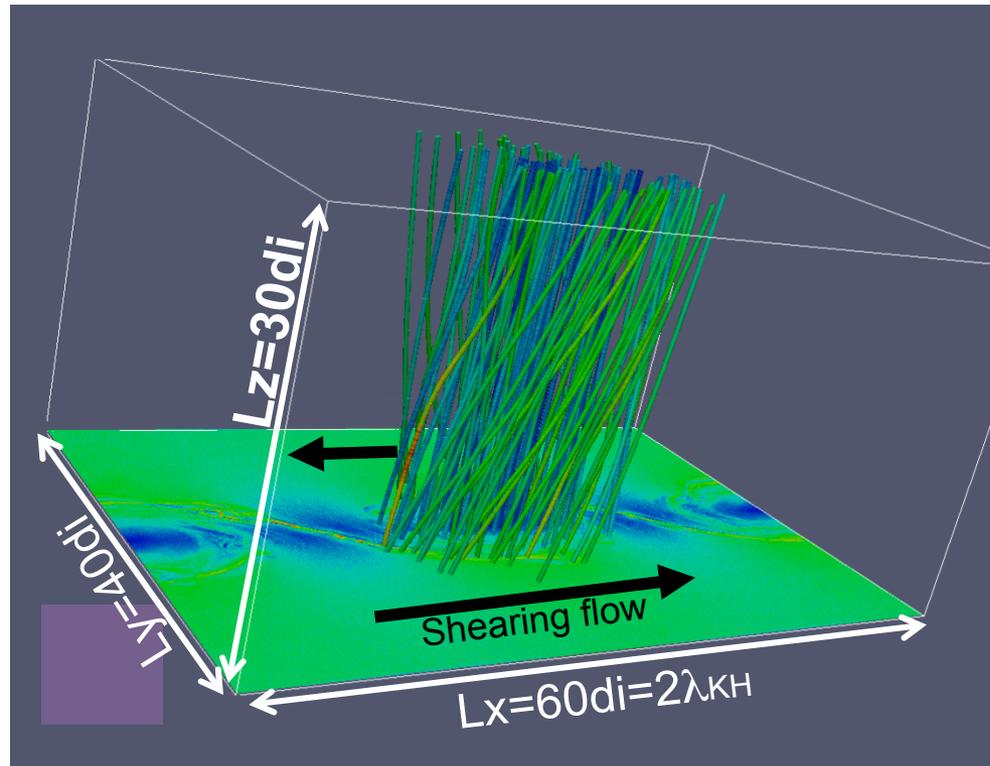
# 2D v.s. 3D (symmetric case)

~ 2D run ~

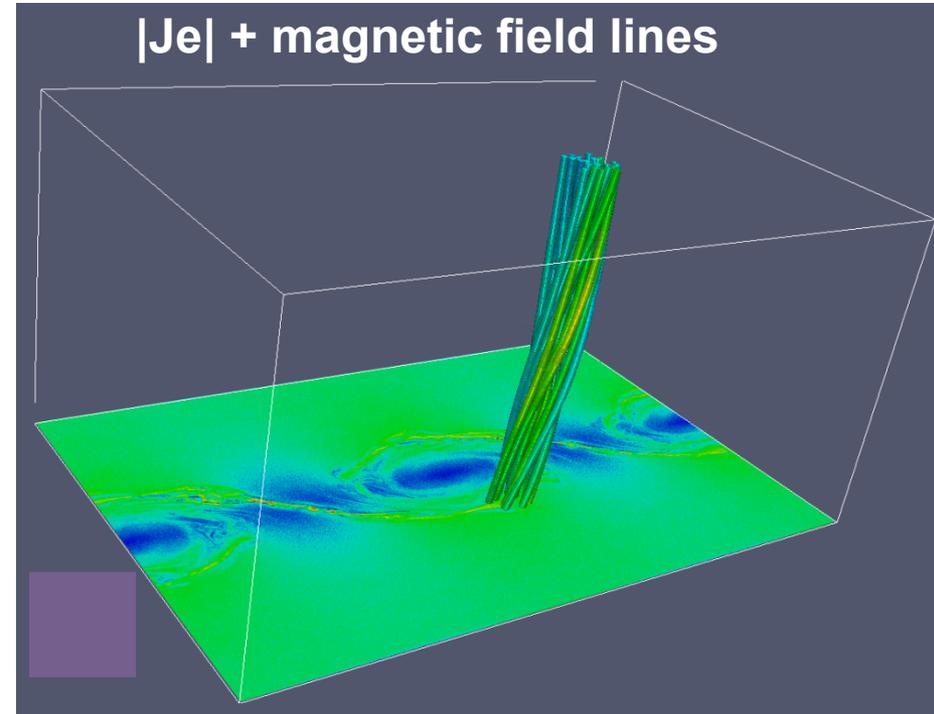
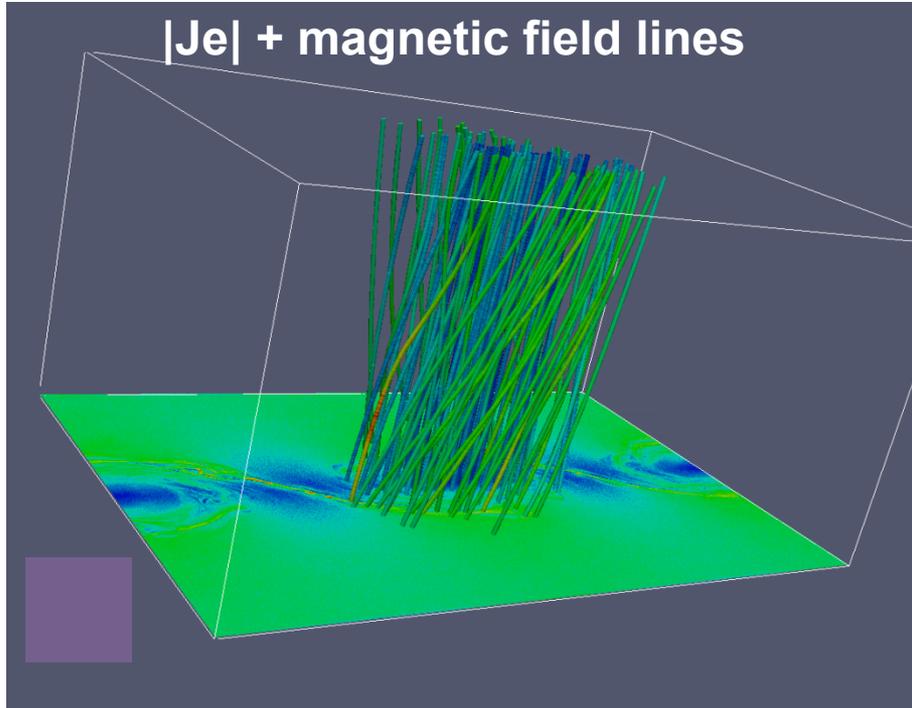
- $L_x \times L_y = 2048 \times 1368 = 60\text{di} \times 40\text{di}$  ( $M_i/M_e = 25$ )
- $7.3 \times 10^8$  particles (260 particles/cell)

~ 3D run ~

- $L_x \times L_y \times L_z = 2048 \times 1368 \times 1024 = 60\text{di} \times 40\text{di} \times 30\text{di}$  ( $M_i/M_e = 25$ )
- $7.5 \times 10^{11}$  particles (260 particles/cell)

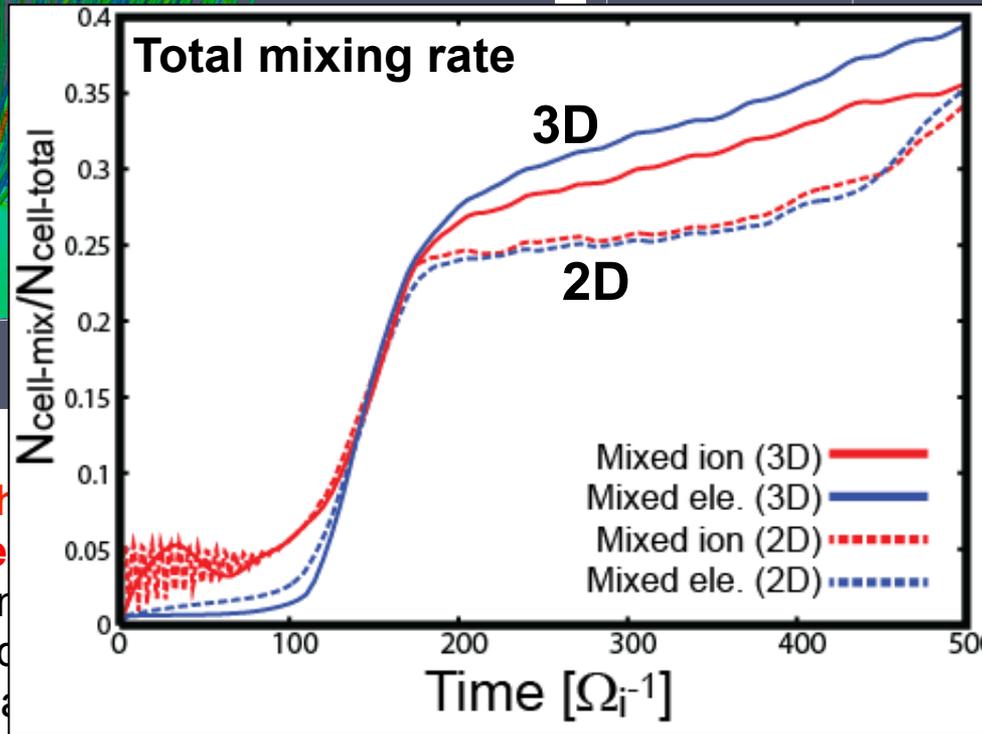
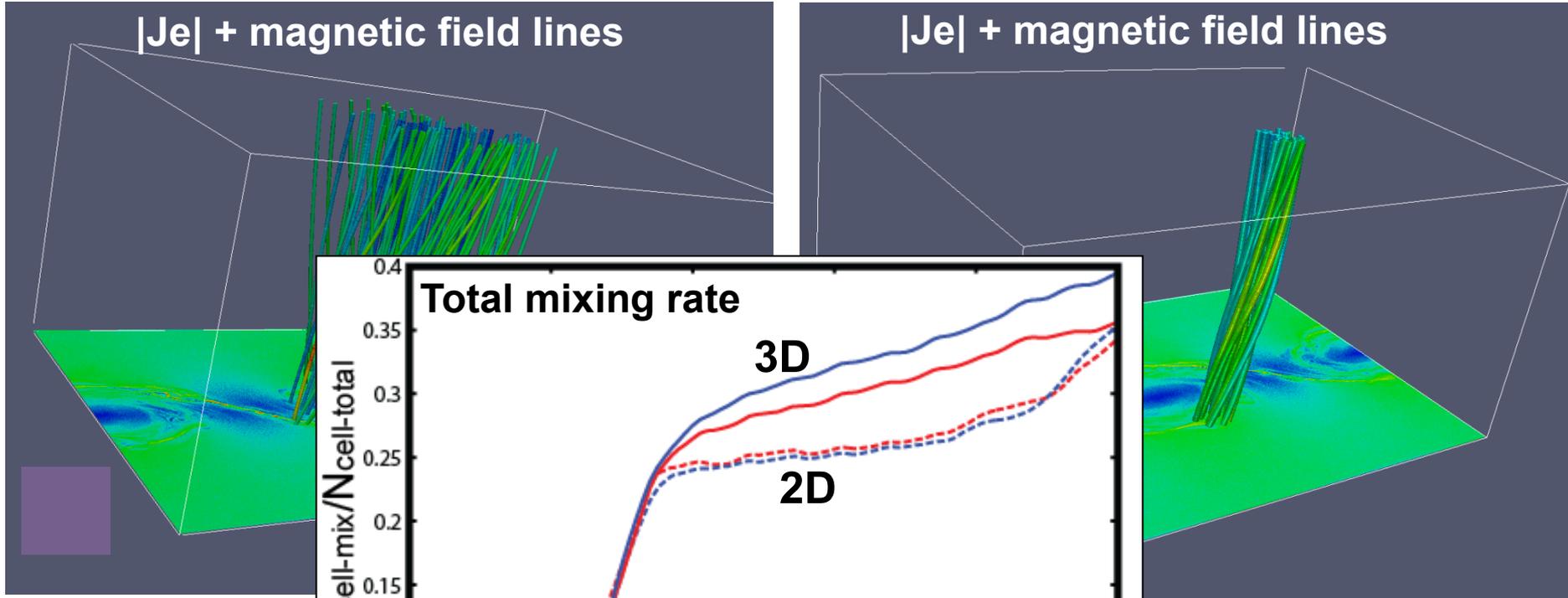


# 2D v.s. 3D (symmetric case)



- ◆ **3D magnetic shear** exists along the edge of the vortex.
  - **3D reconnection** occurs and produces flux ropes there.
  - 3D reconnection disturbs the vortex structure and causes turbulence.
  - Inside and outside of the vortex are connected by 3D reconnection (flux ropes).
    - Plasmas inside the vortex can be transported outside the vortex.
    - Mixed area broadens even outside the vortex.
- ◆ **Consequently, plasma mixing in 3D progresses more quickly than 2D.**

# 2D v.s. 3D (symmetric case)

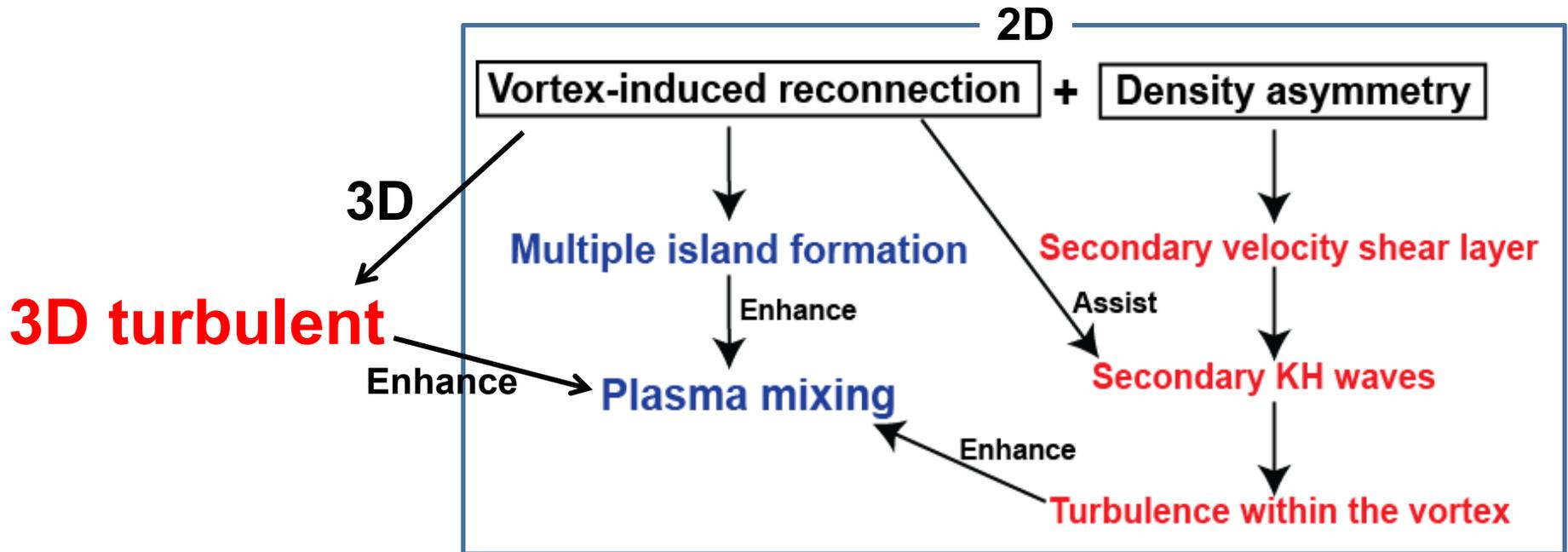


- ◆ 3D magnetic shears
- 3D reconnection
- 3D reconnection
- Inside and outside
- Plasma mixing
- Mixed area broadens even outside the vortex.

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◆ Consequently, plasma mixing in 3D progresses more quickly than 2D.

# Summary



➔ In 3D, the vortex layer becomes turbulent even in the simple symmetric case.

➔ 3D turbulence enhances the plasma mixing.

➔ The KH vortex in collisionless space plasmas would generally cause turbulence and efficient plasma mixing.